

A Family of Well Behaved Uncharged Model of Vaidya–Tikekar Type Super-Dense Star

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Abstract: In the present article a model of well-behaved uncharged super-dense star with surface density 2×10^{14} gm/cm³ is developed by considering a static spherically symmetric metric with $t = \text{const}$ hyper surfaces as hyperboloid. So far well-behaved model described by such metric could not be obtained for uncharged fluid. Maximum mass of the star is found to be $0.343121M_{\odot}$ and the corresponding radius is 9.576299 Km. The red shift at the centre and on the surface are given as 0.076792 and 0.031726 respectively.

Key Words: Exact solutions, Uncharged fluids, Super-dense stars, General relativity.

1. Introduction

In short, the uncharged fluids so far obtained are not well behaved. In the present article authors have developed a uncharged super-dense models of above kind which are well behaved. Well behaved formality of models means the models satisfying the conditions $c^2 \rho \geq p \geq 0$, $\frac{dp}{dr} < 0$, $\frac{d\rho}{dr} < 0$, $\frac{p}{c^2 \rho} < 1$, $\frac{dp}{c^2 d\rho} < 1$, $\frac{d}{dr} \left(\frac{p}{c^2 \rho} \right) < 0$, $\frac{d}{dr} \left(\frac{dp}{c^2 d\rho} \right) < 0$, where ρ and p denote the energy density and pressure of the source material. Also $\frac{dp}{d\rho}$ denotes the velocity of sound in through the model.

2. Field Equations

Let us consider the metric to be

$$(2.1) \quad ds^2 = -e^\lambda dr^2 - r^2 (d\theta^2 + \sin^2 \theta d\phi^2) + e^\nu dt^2,$$

where the functions $\lambda(r)$ and $\nu(r)$ are to satisfy the Einstein field equation for perfect fluid distributions

$$(2.2) \quad R_j^i - \frac{1}{2} R \delta_j^i = -\kappa [(c^2 \rho + p) v^i v_j - p \delta_j^i],$$

where $\kappa = \frac{8\pi G}{c^4}$, ρ , p and v^i denote energy density, fluid pressure and flow vector of the fluid, respectively.

The field equation (2.2) with respect to the metric equation (2.1) give¹

$$(2.3) \quad 8\pi T_1^1 = -\frac{\nu'}{r} e^{-\lambda} + \frac{(1 - e^{-\lambda})}{r^2},$$

$$(2.4) \quad 8\pi T_2^2 = 8\pi T_3^3 = -\left[\frac{\nu''}{2} - \frac{\lambda' \nu'}{4} + \frac{\nu'^2}{4} + \frac{\nu' - \lambda'}{2r} \right] e^{-\lambda},$$

$$(2.5) \quad 8\pi T_4^4 = \frac{\lambda'}{r} e^{-\lambda} + \frac{(1 - e^{-\lambda})}{r^2}.$$

At the pressure free interface $r = a$ the charged fluid sphere is expected to join with the Reissner-Nordstrom metric

$$(2.6) \quad ds^2 = -\left(1 - \frac{2m}{r} + \frac{e^2}{r^2}\right)^{-1} dr^2 - r^2 (d\theta^2 + \sin^2 \theta d\phi^2) + \left(1 - \frac{2m}{r} + \frac{e^2}{r^2}\right) dt^2,$$

where m is the gravitational mass of the distribution such that $m = \mu(a)$, while

$$(2.7) \quad \mu(a) = \frac{\kappa}{2} \int_0^a \rho(r) r^2 dr,$$

$\mu(a)$ is the mass².

The space-time with its hyper surface $t = \text{constant}$ as hyperboloid is characterized by the metric (2.1) with the metric potential g_{11} as ³⁻¹⁰

$$(2.8) \quad e^\lambda = \frac{(K + Cr^2)}{K(1 + Cr^2)}, \quad K > 1,$$

where K is parameter and C being a constant¹¹⁻¹².

Now from the equations (2.3), (2.5) and (2.8), we get

$$(2.9) \quad \frac{(K + Cr^2)}{K(1 + Cr^2)} \left[-\frac{2y'}{ry} + \frac{C(K-1)}{(K + Cr^2)} \right] = -\kappa p,$$

$$(2.10) \quad \frac{C(K-1)(3 + Cr^2)}{K(1 + Cr^2)^2} = \kappa c^2 \rho,$$

where $e^\nu = y^2$ and y satisfies the isotropy condition ($T_1^1 = T_2^2 = T_3^3$)

$$(2.11) \quad (K + Cr^2)(1 + Cr^2)y'' - \frac{1}{r}[(K + Cr^2)(1 + Cr^2) + (K-1)Cr^2]y' + (K-1)C^2r^2y = 0.$$

Now the equation (2.11) can be written as

$$(2.12) \quad (1 - X^2) \frac{d^2 y}{dX^2} + X \frac{dy}{dX} + \alpha y = 0,$$

where $\alpha = (1 - K)$ and

$$(2.13) \quad X = \sqrt{\frac{K}{(K-1)}} \sqrt{1 + \frac{Cr^2}{K}}, \quad K > 1.$$

Here we consider $K > 1$, then $X > 1$ and we adopt a different approach to get a solution, which as follows

$$(2.14) \quad y(X) = (X^2 - 1)Z(X),$$

which sends the equation (2.12) to the equation

$$(2.15) \quad \frac{d^2 Z}{dX^2} + IZ = 0,$$

where

$$(2.16) \quad I = \frac{\alpha}{(1-X^2)} - \frac{(2+3X^2)}{4(1-X^2)^2}.$$

In the present equation (2.15) can be solved easily if set $K = \frac{7}{4}$ consequently the equations (2.13) and (2.16) now read as

$$(2.17) \quad X = \sqrt{\frac{(7+4Cr^2)}{3}},$$

$$(2.18) \quad I = -\frac{5}{4(1-X^2)^2},$$

and the equation (2.15) takes the form

$$(2.19) \quad Z'' - \frac{5}{4(1-X^2)^2} Z = 0,$$

which admits the solution

$$Z(X) = (X+1) \left[C_1 \left| \frac{X+1}{X-1} \right|^{(\beta-0.5)} + C_2 \left| \frac{X+1}{X-1} \right|^{(-\beta-0.5)} \right], \quad \beta = \frac{3}{4}.$$

Now from the equation (2.14), we have

$$(2.20) \quad y(X) = (X^2-1)(X+1) \left[C_1 \left| \frac{X+1}{X-1} \right|^{(\beta-0.5)} + C_2 \left| \frac{X+1}{X-1} \right|^{(-\beta-0.5)} \right],$$

which can be written in term of r using the equation (2.17) as (2.21)

$$y(r) = \frac{4(1+Cr^2)}{3} \left(\sqrt{\frac{(7+4Cr^2)}{3}} + 1 \right) \left[C_1 \left| \frac{\sqrt{\frac{(7+4Cr^2)}{3}} + 1}{\sqrt{\frac{(7+4Cr^2)}{3}} - 1} \right|^{(\beta-0.5)} + C_2 \left| \frac{\sqrt{\frac{(7+4Cr^2)}{3}} + 1}{\sqrt{\frac{(7+4Cr^2)}{3}} - 1} \right|^{(-\beta-0.5)} \right].$$

The corresponding expression for energy density and pressure can be had from the equations (2.9), (2.10), and (2.21) as

$$(2.22) \quad \kappa C^2 \rho = \frac{3C(3 + Cr^2)}{7(1 + Cr^2)^2},$$

$$(2.23) \quad \kappa p = -\frac{(7 + 4Cr^2)}{7(1 + Cr^2)} \left[-\frac{2y'}{ry} + \frac{3C}{(7 + 4Cr^2)} \right].$$

Now $p_{(r=a)} = 0$ gives

$$(2.24) \quad \frac{C_2}{C_1} = \frac{DN \left| \frac{\sqrt{\frac{(7 + 4Ca^2)}{3}} + 1}{\sqrt{\frac{(7 + 4Ca^2)}{3}} - 1} \right|^{(\beta-0.5)} - N_1 N_3 \left| \frac{\sqrt{\frac{(7 + 4Ca^2)}{3}} + 1}{\sqrt{\frac{(7 + 4Ca^2)}{3}} - 1} \right|^{(\beta-0.5)} + D_1}{N_1 N_3 \left| \frac{\sqrt{\frac{(7 + 4Ca^2)}{3}} + 1}{\sqrt{\frac{(7 + 4Ca^2)}{3}} - 1} \right|^{(-\beta-0.5)} - DN \left| \frac{\sqrt{\frac{(7 + 4Ca^2)}{3}} + 1}{\sqrt{\frac{(7 + 4Ca^2)}{3}} - 1} \right|^{(-\beta-0.5)} + D_2},$$

where

$$D = \left(\frac{4(1 + Ca^2)}{3} \right)^{\left(\frac{1}{4}\right)} \left(\sqrt{\frac{(7 + 4Ca^2)}{3}} + 1 \right), \quad N = \frac{3}{(7 + 4Ca^2)}$$

$$N_1 = \left(\frac{4(1 + Ca^2)}{3} \right)^{\left(\frac{1}{4}\right)} + \frac{\sqrt{\frac{(7 + 4Ca^2)}{3}} \left(\sqrt{\frac{(7 + 4Ca^2)}{3}} + 1 \right)}{2 \left(\frac{4(1 + Ca^2)}{3} \right)^{\left(\frac{3}{4}\right)}},$$

$$N_2 = \frac{2 \left(\frac{4(1 + Ca^2)}{3} \right)^{\left(\frac{1}{4}\right)} \left(\sqrt{\frac{(7 + 4Ca^2)}{3}} + 1 \right)}{\left(\sqrt{\frac{(7 + 4Ca^2)}{3}} - 1 \right)^2},$$

$$N_3 = \frac{8(7 + 4Ca^2)}{21(1 + Ca^2) \sqrt{\frac{7 + 4Ca^2}{3}}},$$

$$D_1 = N_2 N_3 (\beta - 1) \left| \frac{\sqrt{\frac{(7 + 4Ca^2)}{3}} + 1}{\sqrt{\frac{(7 + 4Ca^2)}{3}} - 1} \right|^{(\beta - 1.5)},$$

$$D_2 = N_2 N_3 (\beta + 1) \left| \frac{\sqrt{\frac{(7 + 4Ca^2)}{3}} + 1}{\sqrt{\frac{(7 + 4Ca^2)}{3}} - 1} \right|^{(-\beta - 1.5)}.$$

The expression for the density and pressure gradient can be written as

$$(2.25) \quad \kappa C^2 \frac{d\rho}{dr} = C^2 r \left[-\frac{6(5 + Cr^2)}{7(1 + Cr^2)^3} \right],$$

(2.26)

$$\kappa \frac{dp}{dr} = C^2 r \left[\frac{\left[y(r) \left\{ \frac{N_4 \left(M_5 + \frac{C_2}{C_1} M_6 \right)}{+ M_4 N_5} \right\} - \frac{N_7 \left(M_8 - \frac{C_2}{C_1} M_9 \right)}{+ M_7 N_6} \right] \frac{8(7 + 4Cr^2)}{21(1 + Cr^2) \sqrt{\frac{(7 + 4Cr^2)}{3}}} + \left[M_{10} \{ N_4 N_5 - N_7 N_6 \} - \left[(N_4 N_5 - N_7 N_6)^2 \frac{32}{21(1 + Cr^2)} \right] \right]}{(y(r))^2} + \frac{24}{(7 + 4Cr^2)^2} \right],$$

where

$$N_4 = \left(\frac{4(1 + Cr^2)}{3} \right)^{\left(\frac{1}{4}\right)} + \frac{\sqrt{\frac{(7 + 4Cr^2)}{3}} \left(\sqrt{\frac{(7 + 4Cr^2)}{3}} + 1 \right)}{2 \left(\frac{4(1 + Cr^2)}{3} \right)^{\left(\frac{3}{4}\right)}},$$

$$N_5 = \left[(\beta - 0.5) \left| \frac{\sqrt{\frac{(7+4Cr^2)}{3}} + 1}{\sqrt{\frac{(7+4Cr^2)}{3}} - 1} \right|^{(\beta-0.5)} + (\beta + 0.5) \frac{C_2}{C_1} \left| \frac{\sqrt{\frac{(7+4Cr^2)}{3}} + 1}{\sqrt{\frac{(7+4Cr^2)}{3}} - 1} \right|^{(-\beta-0.5)} \right],$$

$$N_6 = \left[(\beta - 0.5) \left| \frac{\sqrt{\frac{(7+4Cr^2)}{3}} + 1}{\sqrt{\frac{(7+4Cr^2)}{3}} - 1} \right|^{(\beta-1.5)} - (\beta + 0.5) \frac{C_2}{C_1} \left| \frac{\sqrt{\frac{(7+4Cr^2)}{3}} + 1}{\sqrt{\frac{(7+4Cr^2)}{3}} - 1} \right|^{(-\beta-1.5)} \right],$$

$$N_7 = \frac{2 \left(\frac{4(1+Cr^2)}{3} \right)^{\left(\frac{1}{4}\right)} \left(\sqrt{\frac{(7+4Cr^2)}{3}} + 1 \right)}{\left(\sqrt{\frac{(7+4Cr^2)}{3}} - 1 \right)^2},$$

$$M_4 = \frac{1}{2} \frac{\left[\left(\frac{4(1+Cr^2)}{3} \right)^{\left(\frac{3}{4}\right)} \left[\frac{4}{3} + \frac{4}{3} \left(\sqrt{\frac{(7+4Cr^2)}{3}} + 1 \right) \right] - 2 \left[\left(\frac{4(1+Cr^2)}{3} \right)^{\frac{-1}{4}} \left(\sqrt{\frac{(7+4Cr^2)}{3}} + 1 \right) \right] \right]}{\left(\frac{4(1+Cr^2)}{3} \right)^{\left(\frac{3}{2}\right)} \left[\left(\frac{(7+4Cr^2)}{3} \right)^{\frac{-1}{2}} \right] \left[\sqrt{\frac{(7+4Cr^2)}{3}} \right]} + \frac{2}{3} \left(\frac{4(1+Cr^2)}{3} \right)^{\frac{-3}{4}},$$

$$M_5 = \frac{-8(\beta - 0.5)}{3 \sqrt{\frac{(7+4Cr^2)}{3}} \left(\sqrt{\frac{(7+4Cr^2)}{3}} - 1 \right)^2} \left| \frac{\sqrt{\frac{(7+4Cr^2)}{3}} + 1}{\sqrt{\frac{(7+4Cr^2)}{3}} - 1} \right|^{(\beta-1.5)},$$

$$M_6 = \frac{8(\beta + 0.5)}{3\sqrt{\frac{(7+4Cr^2)}{3}}\left(\sqrt{\frac{(7+4Cr^2)}{3}}-1\right)^2} \left| \frac{\sqrt{\frac{(7+4Cr^2)}{3}}+1}{\sqrt{\frac{(7+4Cr^2)}{3}}-1} \right|^{(-\beta-1.5)},$$

$$M_7 = \frac{2 \left[\left(\sqrt{\frac{(7+4Cr^2)}{3}}-1 \right)^2 \left[\frac{2}{3} \left(\frac{4(1+Cr^2)}{3} \right)^{\frac{-3}{4}} \left(\sqrt{\frac{(7+4Cr^2)}{3}}+1 \right) + \frac{4}{3} \left(\frac{4(1+Cr^2)}{3} \right)^{\frac{1}{4}} \right] \right.}{\left[\left(\frac{(7+4Cr^2)}{3} \right)^{\frac{-1}{2}} \right]} - \left[\frac{8}{3} \left(\sqrt{\frac{(7+4Cr^2)}{3}}+1 \right) \left(\frac{4(1+Cr^2)}{3} \right)^{\frac{1}{4}} \left(\sqrt{\frac{(7+4Cr^2)}{3}}-1 \right) \left(\frac{(7+4Cr^2)}{3} \right)^{\frac{-1}{2}} \right]} \right] \left(\sqrt{\frac{(7+4Cr^2)}{3}}-1 \right)^4},$$

$$M_8 = \frac{-8(\beta - 0.5)(\beta - 1.5)}{3\sqrt{\frac{(7+4Cr^2)}{3}}\left(\sqrt{\frac{(7+4Cr^2)}{3}}-1\right)^2} \left| \frac{\sqrt{\frac{(7+4Cr^2)}{3}}+1}{\sqrt{\frac{(7+4Cr^2)}{3}}-1} \right|^{(\beta-2.5)},$$

$$M_9 = \frac{8(\beta + 0.5)(\beta + 1.5)}{3\sqrt{\frac{(7+4Cr^2)}{3}}\left(\sqrt{\frac{(7+4Cr^2)}{3}}-1\right)^2} \left| \frac{\sqrt{\frac{(7+4Cr^2)}{3}}+1}{\sqrt{\frac{(7+4Cr^2)}{3}}-1} \right|^{(-\beta-2.5)},$$

$$M_{10} = \frac{\frac{8}{7} \left[8(1+Cr^2) \sqrt{\frac{(7+4Cr^2)}{3}} - (7+4Cr^2) \right] \left\{ \frac{4}{3} \frac{(1+Cr^2)}{\sqrt{\frac{(7+4Cr^2)}{3}}} \right\} + 2\sqrt{\frac{(7+4Cr^2)}{3}}}{(1+Cr^2)^2(7+4Cr^2)},$$

The expression for velocity of sound can be had from the equations (2.25) and (2.26) as

$$(2.27) \quad \frac{dp}{c^2 d\rho} = \frac{dp/dr}{c^2 d\rho/dr}$$

and the expression of mass

$$(2.28) \quad m(r) = \frac{r}{2} \left[1 - \frac{(7 + 4Cr^2)}{7(1 + Cr^2)} \right],$$

such that

$$e^{-\lambda} = 1 - \frac{2m}{r} + \frac{q^2}{r^2}.$$

3. Physical Conditions to be Satisfied

The physical validity of the charged fluid sphere (CFS) depends upon the following conditions (called reality conditions or energy conditions) inside and on the sphere $r = a$ such that

- (i) $\rho > 0$, $0 \leq r \leq a$,
- (ii) $p > 0$, $r < a$,
- (iii) $p = 0$, $r = a$,
- (iv) $dp/dr < 0$, $d\rho/dr < 0$, $0 < r < a$
- (v) $c^2 \rho \geq p$ weak energy condition (WEC) or $c^2 \rho \geq 3p$ strong energy condition (SEC) $0 \leq r \leq a$.
- (vi) The velocity of sound $(dp/d\rho)^{1/2}$ should be less than that of light throughout the CFS ($0 \leq r \leq a$).
- (vii) $\frac{d}{dr} \left(\frac{p}{c^2 \rho} \right) < 0$.
- (viii) $\frac{d}{dr} \left(\frac{dp}{c^2 d\rho} \right) < 0$.

$$(ix) \quad \text{The adiabatic constant } \gamma = \left(\left(\frac{c^2 \rho + p}{p} \right) \left(\frac{dp}{c^2 d\rho} \right) \right) > 1.$$

Beside the above the smooth joining with the Reissner- Nordström metric, requires the continuity of e^λ , e^ν and q across the pressure free interface $r = a$ and we get

$$(3.1) \quad \frac{7(1 + Ca^2)}{(7 + 4Ca^2)} = 1 - \frac{2m(a)}{a} + \frac{e^2}{a^2},$$

$$(3.2) \quad y^2 = 1 - \frac{2m(a)}{a} + \frac{e^2}{a^2},$$

$$(3.3) \quad p_{(r=a)} = 0.$$

The condition (3.1) is automatically satisfied due to the preposition (2.6) however (3.2) and (3.3) can provide the unique values of arbitrary constants C_1 and C_2 .

4. Conclusion

So, we have been successful in obtaining the uncharged analogues of neutral fluid by considering a spherically symmetric metric with $t = \text{const}$ as hyperboloid $K = \left(\frac{7}{4} \right)$ which is well behaved for $0 < Ca^2 \leq 0.3259$.

Maximum mass and the corresponding radius are found to be $0.343121 M_\odot$ and 9.576299 Km for $Ca^2 = 0.3259$. The red shift at the centre and on the surface are given as 0.076792 and 0.037879 respectively. The adiabatic constant(γ) always more then $4/3$, which is sufficient condition for the stability of the model. Heintzmann and Hillebrandt¹³ proposed that neutron star models with anisotropic equation of state are stable if $\gamma > 4/3$. In the absence of charge the model reduces to the model obtained by Gupta and Jasim³.

$$D = \frac{8\pi G}{c^2} a^2 \rho, \quad P = \frac{8\pi G}{c^4} a^2 p,$$

$$c = 2.997 \times 10^{10} \text{ cm/s}, G = 6.673 \times 10^{-8} \text{ cm}^3 / \text{gs}^2, M_\odot = 1.475 \text{ km}$$

where γ denotes the adiabatic constant and it is given by the expression $\gamma = \left(\left(\frac{c^2 \rho + p}{p} \right) \left(\frac{dp}{c^2 d\rho} \right) \right)$. z_0 and z_a are red shift at the centre and surface $r = a$ respectively. q stands for charge.

Table1: $Ca^2 = 0.007$

Radius = 1.550058Km, $M = 0.001571M_\odot$, $z_0 = 0.001767$, $z_a = 0.001015$							
X	(P)	(D)	(D-3P)	(q)	$dp/c^2 d\rho$	P/D	γ
0	0.000016	0.009	0.008953	0	0.150447	0.001737	86.74692
0.2	0.000015	0.008998	0.008953	0	0.150442	0.001667	90.37198
0.4	0.000013	0.008993	0.008954	0	0.150424	0.001458	103.3186
0.6	0.00001	0.008985	0.008955	0	0.150396	0.00111	135.6846
0.8	0.000006	0.008973	0.008956	0	0.150355	0.000623	241.4125
1	0	0.008958	0.008958	0	0.150301	0	Inf

Table 2: $Ca^2 = 0.01$

Radius = 1.850827Km, $M = 0.002671M_\odot$, $z_0 = 0.002523$, $z_a = 0.001446$							
x	(P)	(D)	(D-3P)	(q)	$dp/c^2 d\rho$	P/D	γ
0	0.000032	0.012857	0.012762	0	0.150637	0.002474	61.03306
0.2	0.000031	0.012854	0.012762	0	0.150629	0.002375	63.58682
0.4	0.000027	0.012843	0.012763	0	0.150605	0.002076	72.70729
0.6	0.00002	0.012826	0.012766	0	0.150564	0.001579	95.50795
0.8	0.000011	0.012802	0.012768	0	0.150506	0.000886	169.9882
1	0	0.012772	0.012772	0	0.150429	0	Inf

Table 3: $Ca^2 = 0.1$

<i>Radius = 5.683255Km, $M = 0.075315M_{\odot}$, $z_0 = 0.024715$, $z_a = 0.013521$</i>							
x	(P)	(D)	(D−3P)	(q)	$dp/c^2d\rho$	P/D	γ
0	0.00291	0.128571	0.119842	0	0.155919	0.022633	7.045012
0.2	0.002777	0.128229	0.119898	0	0.155861	0.021657	7.352728
0.4	0.002387	0.127211	0.120049	0	0.155674	0.018768	8.450453
0.6	0.001767	0.125541	0.12024	0	0.155314	0.014076	11.18929
0.8	0.000955	0.123261	0.120395	0	0.154709	0.00775	20.11606
1	0	0.120425	0.120425	0	0.153769	0	Inf

Table 4: $Ca^2 = 0.2$

<i>Radius = 7.788896Km, $M = 0.189235M_{\odot}$, $z_0 = 0.048364$, $z_a = 0.025218$</i>							
x	(P)	(D)	(D−3P)	(q)	$dp/c^2d\rho$	P/D	γ
0	0.010624	0.257143	0.225271	0	0.160953	0.041315	4.056675
0.2	0.010078	0.255777	0.225542	0	0.160883	0.039403	4.243922
0.4	0.008514	0.251745	0.226204	0	0.160614	0.033818	4.909944
0.6	0.006127	0.245242	0.226862	0	0.159975	0.024983	6.563333
0.8	0.003192	0.236586	0.22701	0	0.158709	0.013492	11.92217
1	0	0.22619	0.22619	0	0.156507	0	Inf

Table 5: $Ca^2 = 0.3$

<i>Radius = 9.257469Km, $M = 0.31142M_{\odot}$, $z_0 = 0.071058$, $z_a = 0.035438$</i>							
x	(P)	(D)	(D-3P)	(q)	$dp/c^2d\rho$	P/D	γ
0	0.021969	0.385714	0.319807	0	0.165272	0.056957	3.066963
0.2	0.020715	0.382647	0.320501	0	0.165225	0.054137	3.217194
0.4	0.017196	0.373667	0.32208	0	0.164952	0.046019	3.749394
0.6	0.012042	0.359428	0.323301	0	0.164082	0.033504	5.061399
0.8	0.00606	0.34096	0.322779	0	0.162086	0.017774	9.281191
1	0	0.319527	0.319527	0	0.158377	0	Inf

Table 6: $Ca^2 = 0.3259$

<i>Radius=9.257469Km, $M=0.31142M_{\odot}$, $z_0=0.071058$, $z_a=0.035438$</i>							
x	(P)	(D)	(D-3P)	(q)	$dp/c^2d\rho$	P/D	γ
0	0.025391	0.419014	0.342841	0	0.16629	0.060597	2.910492
0.2	0.023904	0.415397	0.343683	0	0.166254	0.057546	3.055321
0.4	0.019754	0.404827	0.345567	0	0.16599	0.048795	3.567786
0.6	0.013738	0.388142	0.346927	0	0.165065	0.035395	4.82859
0.8	0.006854	0.366645	0.346082	0	0.162868	0.018695	8.874689
1	0	0.341915	0.341915	0	0.158737	0	Inf

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